



中国科学技术大学
University of Science and Technology of China



火灾科学国家重点实验室
State Key Laboratory of Fire Science

Effect of wall friction on shock-flame interactions in a hydrogen-air mixture

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2023/10/31

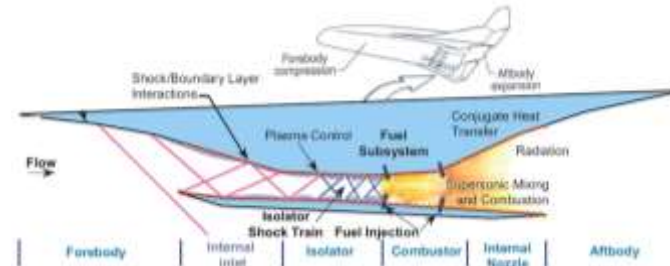
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Shock-Flame Interactions and Instability

Shock Flame interactions (SFI)

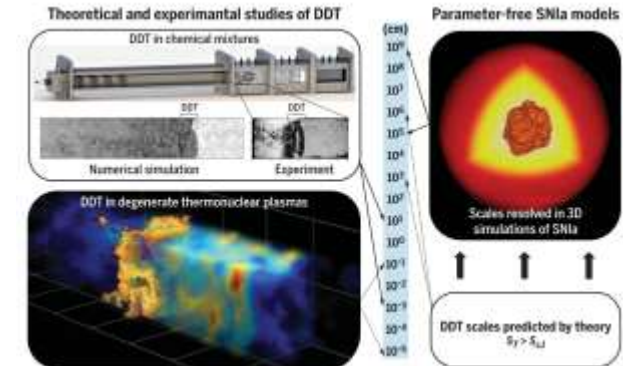
- Extend flame surface
- Compression
- Richtmyer-Meshkov (RM) Instability
- The most important process before flame acceleration and DDT



Hypersonic airbreathing engine



Power Plant Hydrogen Explosion
Muskingum, Ohio, 2007



Supernova Explosions

Poludnenko et al., 2019, Science



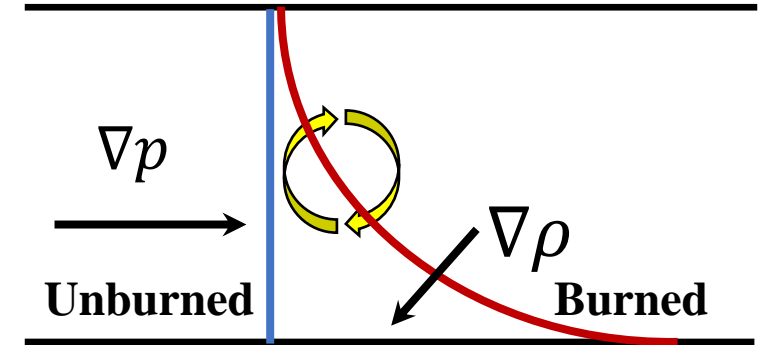
Hydrogen Refueling Station
Explosion Norway, 2019



Motivation

Shock-flame interactions

- Single flame interacts with shock
Flame instability (Richtmyer-Meshkov instability)
Mechanism: Baroclinic torque
Inert RMI model: Impulsive Model, Zhang & Sohn, Meyer and Blewett, Mikaelian...
Reactive RMI model [Yang & Radulescu, 2021, JFM]
- Complicated shock-flame interactions
Shock interacts with multi flames [Bakalis 2021]
Shock interacts with bubble flame [Haehn 2023, CNF] [Diegelmann 2021, CNF] [Thomas 2007 CTM]

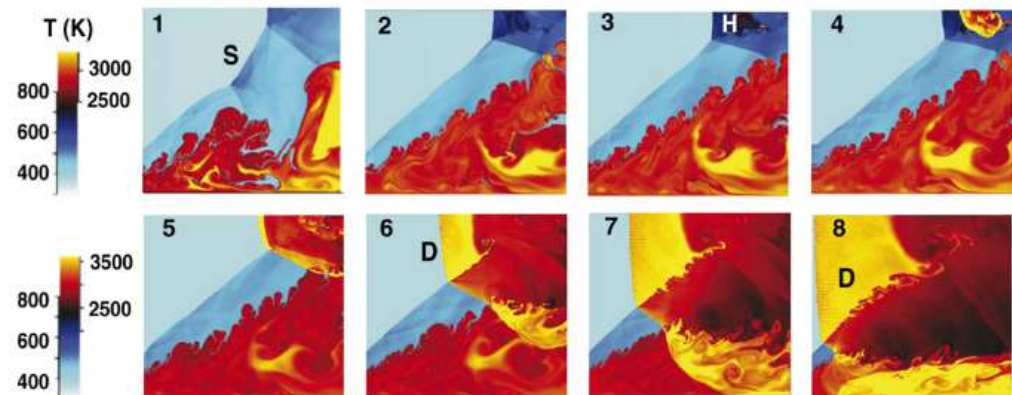




Motivation

Shock-flame interactions

- Reflected shock, flame and boundary later interactions
Shock bifurcation, lambda shock and promotes DDT
[Gamezo 2001 CNF] [Gamezo 2005 PCI] [Yhuel 2023 PCI]



What if shock-flame interactions occur with the existence of wall friction?

Methodology

Physical Model

➤ Initially perturbed flame

Unburned region $x < (x_f + L_y) - \sqrt{L_y^2 - y^2}$

Burned region $x > (x_f + L_y) - \sqrt{L_y^2 - y^2}$

Flame front $x = (x_f + L_y) - \sqrt{L_y^2 - y^2}$

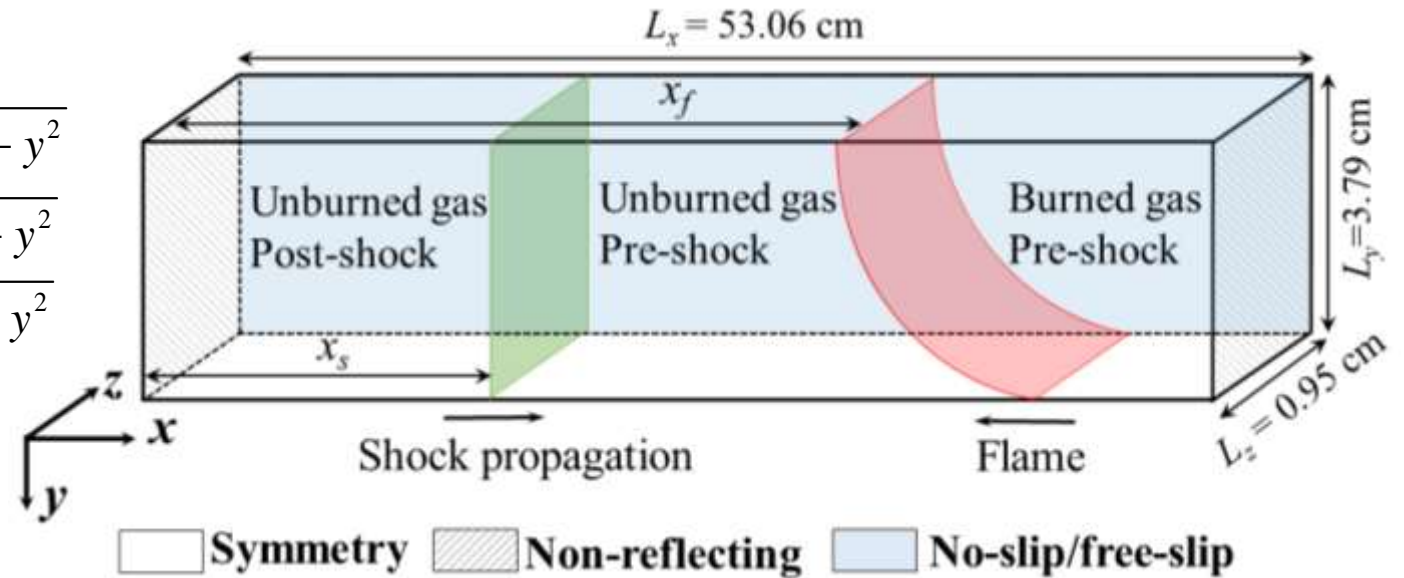
➤ Shock wave

(Rankine-Hugoniot relations)

$$\frac{\rho_2}{\rho_1} = \left(1 + \frac{\gamma + 1}{\gamma - 1} \frac{p_2}{p_1}\right) / \left(\frac{\gamma + 1}{\gamma - 1} + \frac{p_2}{p_1}\right)$$

$$\frac{p_2}{p_1} = \frac{2\gamma}{\gamma + 1} M^2 - \frac{\gamma - 1}{\gamma + 1}$$

.....



Computational domain



Methodology

- **Governing Equations**
 - Reactive, Navier-Stokes Equation
 - Ideal Gas Model
 - Chemical-Diffusive Model
- **Numerical Method**
 - 5th WENO Scheme, HLLC fluxes
 - 3rd Runge-Kutta

$$\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_i u_j) = -\frac{\partial p}{\partial x_i} + \frac{\partial \tau_{ij}}{\partial x_j}$$

$$\frac{\partial}{\partial t}(\rho Y) + \frac{\partial}{\partial x_i}(\rho Y u_i) = \frac{\partial}{\partial t}(\rho D \frac{\partial Y}{\partial x_i}) + \rho \dot{\omega}$$

$$\frac{\partial(\rho E)}{\partial t} + \frac{\partial(\rho u_j E + p u_j)}{\partial x_j} = \frac{\partial}{\partial x_j} \left(K \frac{\partial T}{\partial x_j} \right) + \frac{\partial(\tau_{ij} u_i)}{\partial x_j} - \rho q \dot{\omega}$$

$$p = \rho R T / M$$

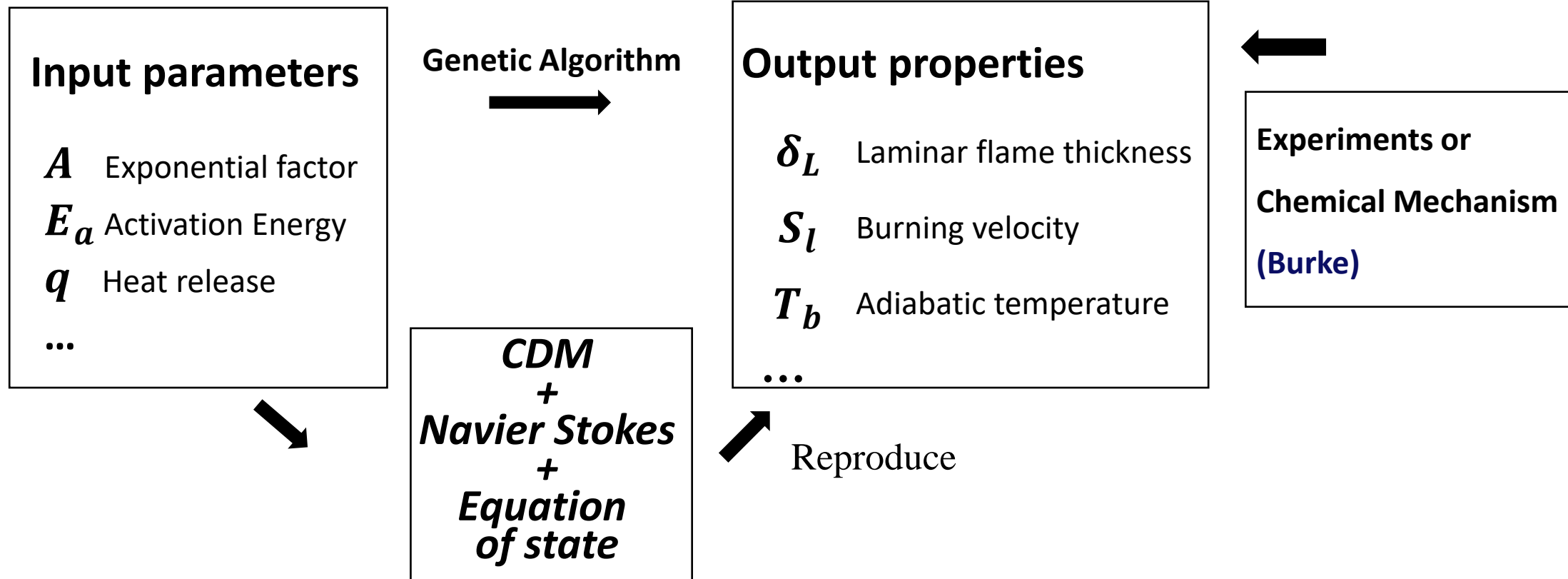
$$E = \frac{p}{\rho(\gamma - 1)} + \frac{1}{2} u_i^2$$

$$\nu = \frac{\nu_0 T^{0.7}}{\rho}, \quad D = \frac{D_0 T^{0.7}}{\rho}, \quad \alpha = \frac{K}{\rho C_p} = \frac{\kappa_0 T^{0.7}}{\rho}$$



Methodology

Arrhenius equation : $\dot{\omega} = \frac{dY}{dt} = A\rho Y \exp(-E_a/RT)$ Y unburned mass fraction



[1] BURKE M P, CHAOS M, JU Y, et al. Comprehensive H₂/O₂ kinetic model for high-pressure combustion [J]. Int J Chem Kinet, 2012, 44(7): 444-74.

[2] KAPLAN C R, ÖZGEN A, ORAN E S. Chemical-diffusive models for flame acceleration and transition-to-detonation: genetic algorithm and optimisation procedure [J]. Combustion Theory and Modelling, 2019, 23(1): 67-86.



Methodology

➤ **Mixture: $\varphi = 1$ H₂-Air, 17.24 kPa, 293K** [Burke, Int J Chem Kinet, 2012]

➤ **Input parameters**

γ	1.1648	Specific heat ratio
M	24.2	Molecular weight
A	1.332×10^8 m ³ /kg-s	Pre-exponential factor
E_a	$33.24 RT_0$	Activation energy
q	$48.70 RT_0/M$	Heat release
κ_0	3.648×10^{-6} kg/s-m-K ^{0.7}	Transport constants

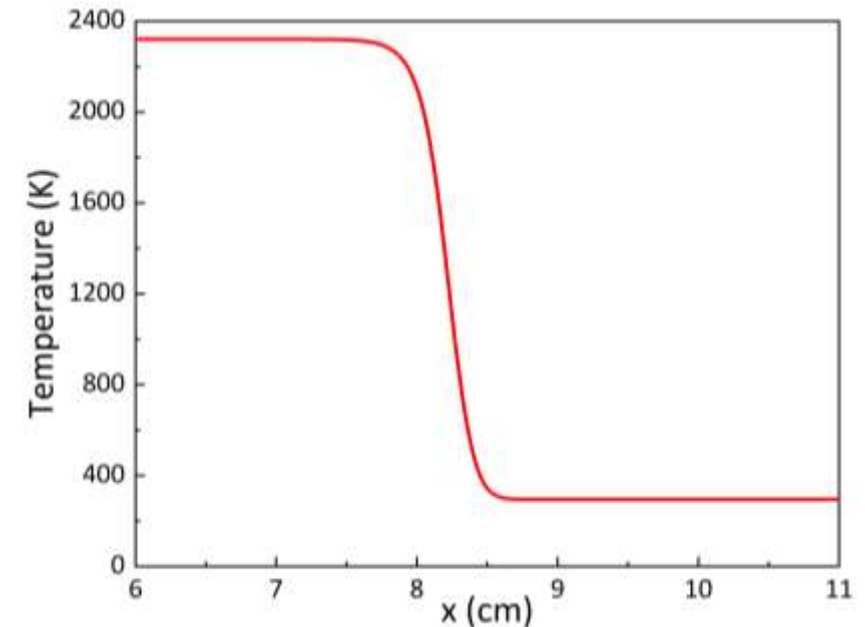
➤ **Output properties**

Laminar burning velocity **1.98 m/s**

Adiabatic flame temperature **2320 K**

Laminar flame thickness **0.375 cm**

➤ **$Pr = 0.1, Le = 1$**



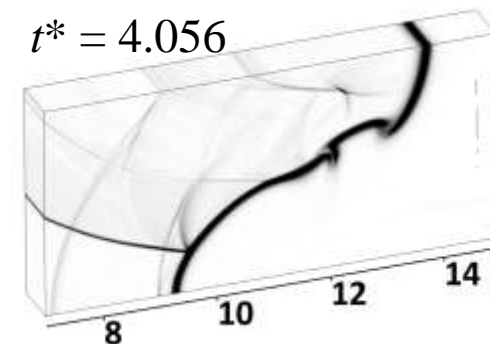
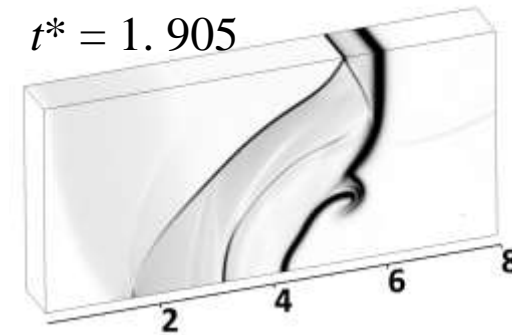
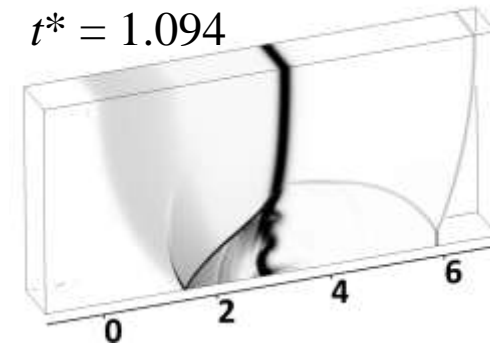
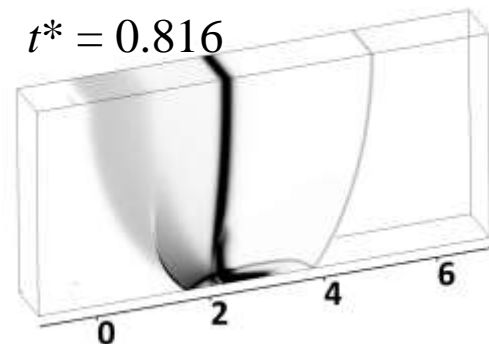
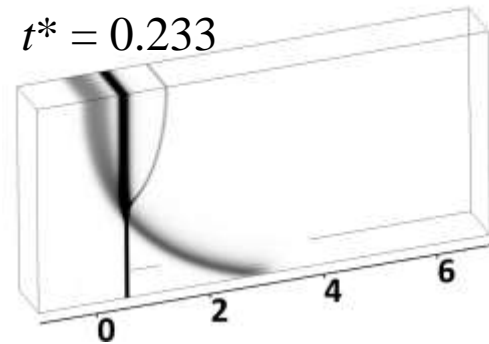
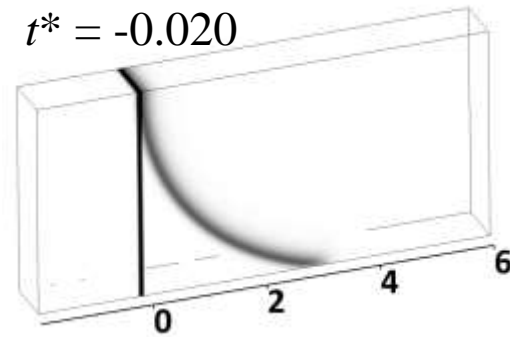
One-dimensional steady flame



Shock-Flame Interaction and Flame Evolution

Free-slip case

- 2D
- Long Neck

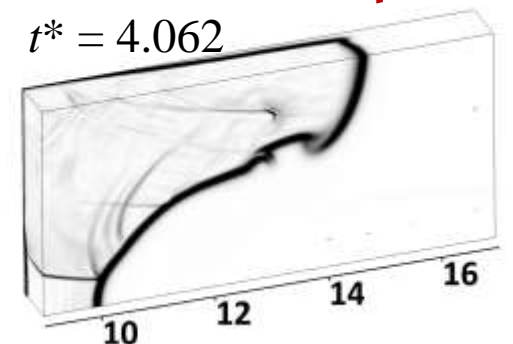
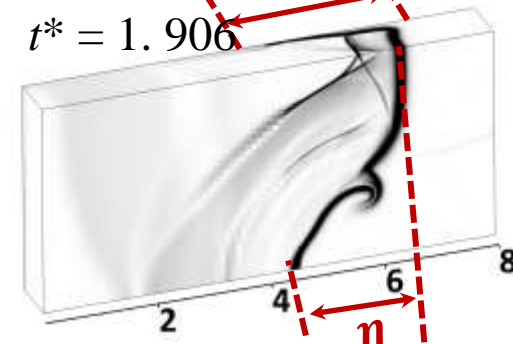
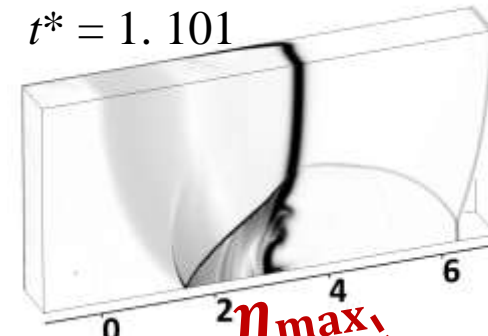
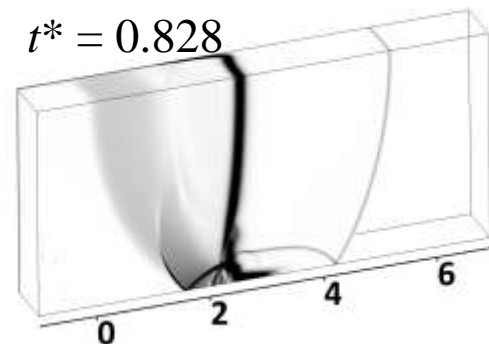
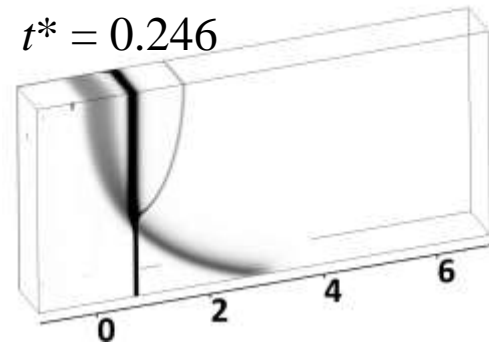
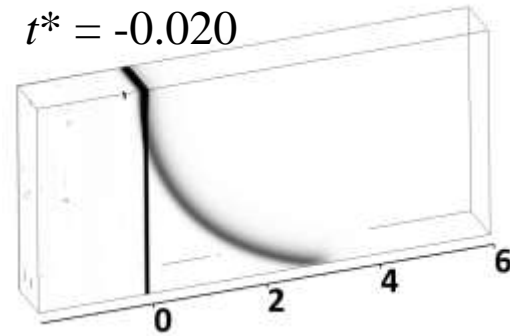




Shock-Flame Interaction and Flame Evolution

No-slip case

- 3D
- Wall friction breaks the 2D
- Flame stretch



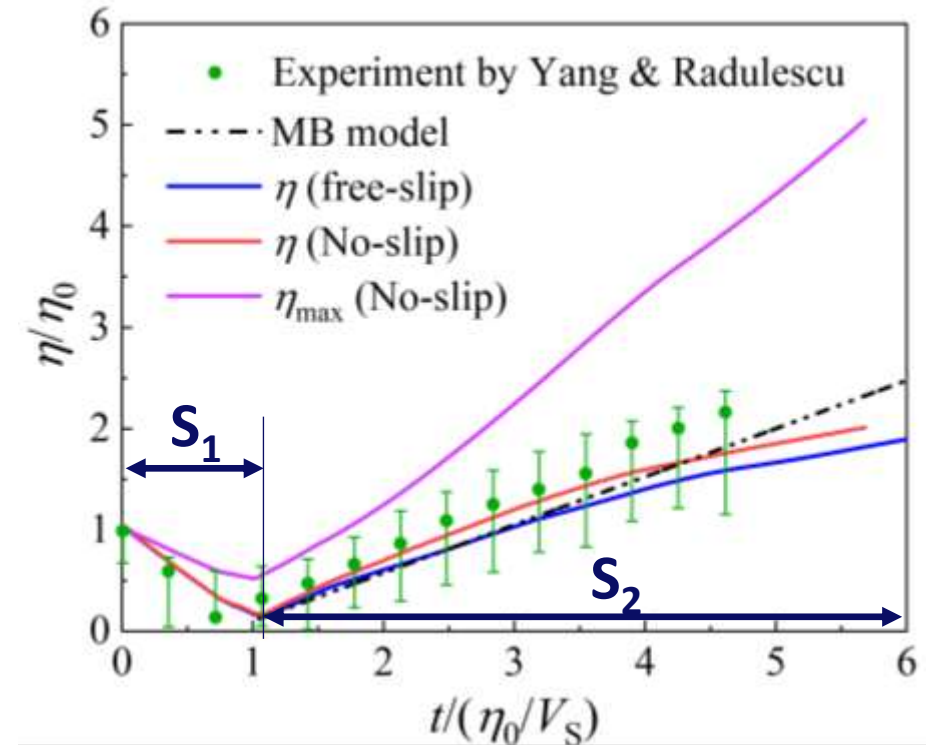


Flame Instability

Inert RMI model

Meyer & Blewett (MB) Model

- Good agreement with experiment [Yang & Radulescu, 2021, JFM]
- Two stages
 - S_1 , Shock compression
 - S_2 , Perturbation growth stages
- Wall friction promotes flame instability

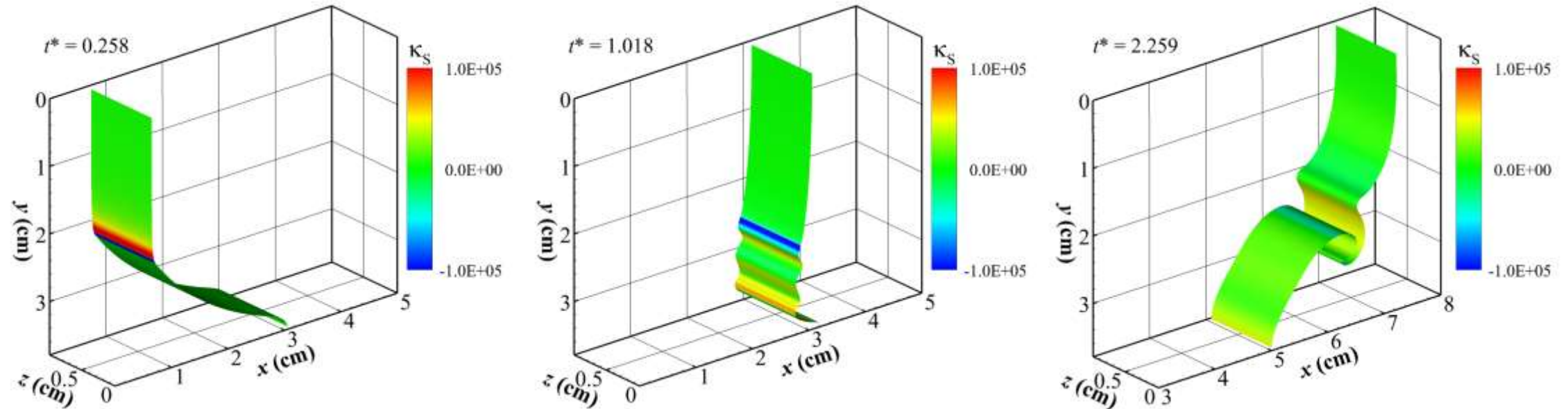


Flame instabilities as a function time



Flame stretch during SFI

Free-slip case



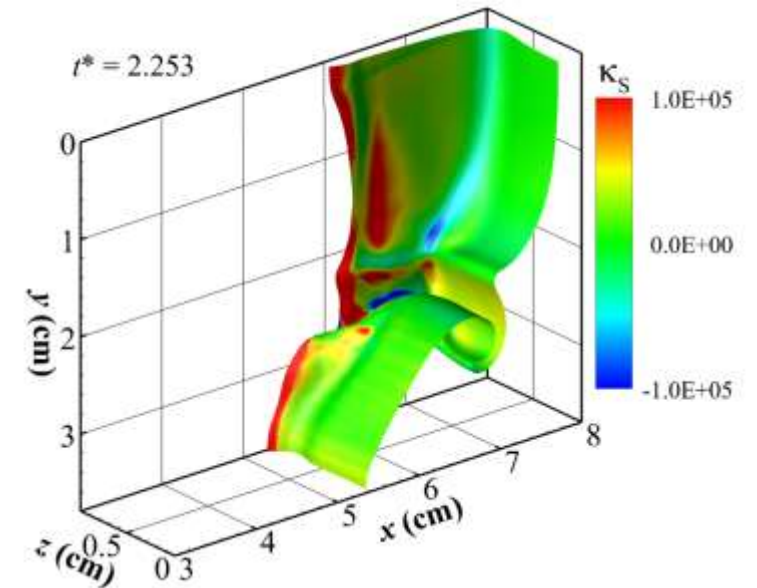
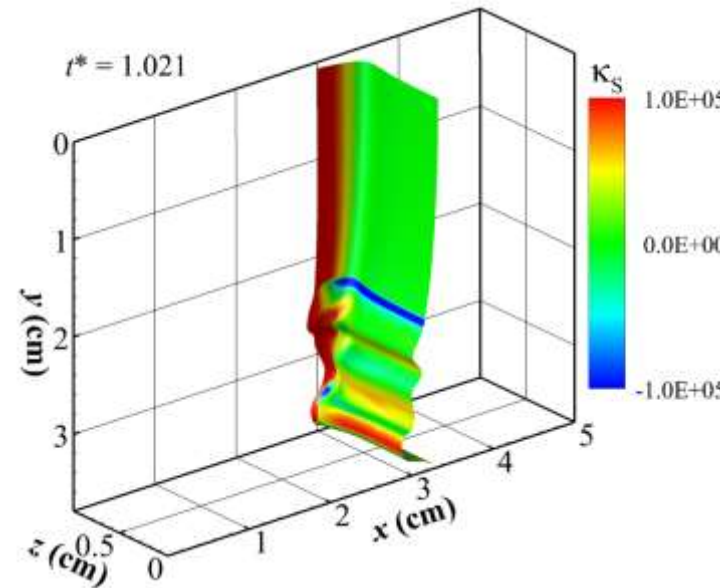
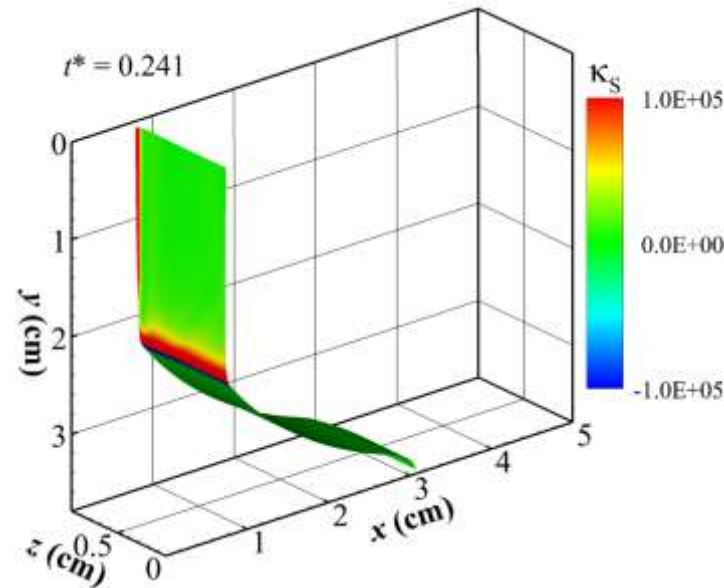
Flame stretch rate

Peak stretch rate appears at the intersection of shock and flame
No obvious stretch after shock passage



Flame Stretch during SFI

No-slip case



Flame stretch rate

Heavy, and continuous stretch near the wall (10^5)

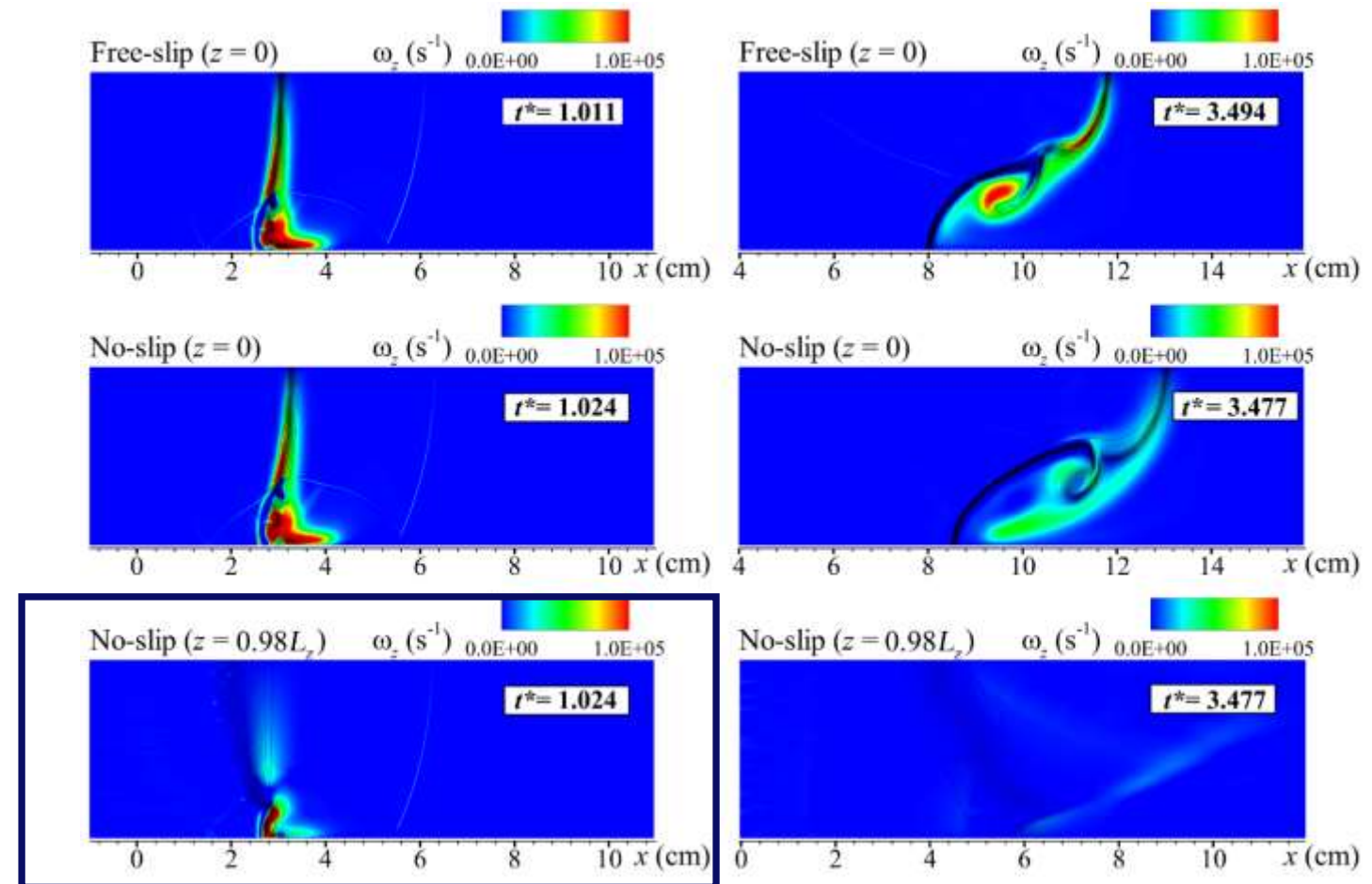
Thin boundary layer

Larger flame surface area



Vortex Dynamics

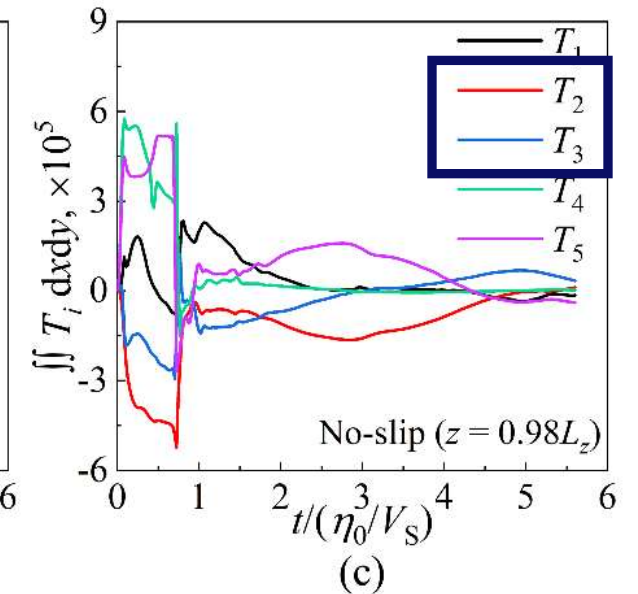
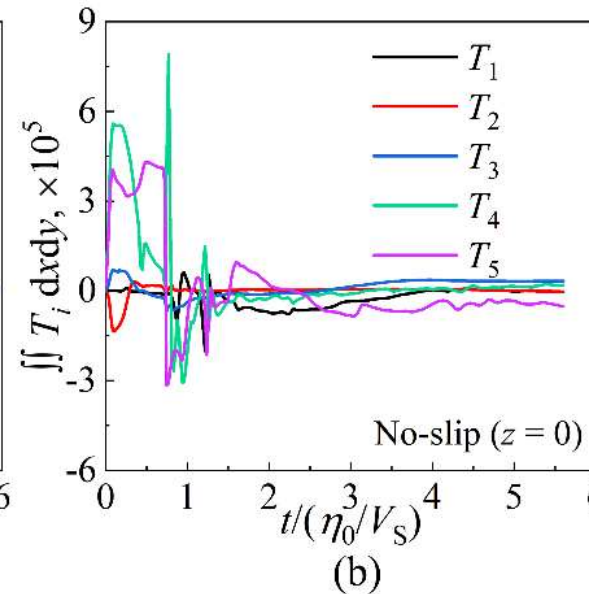
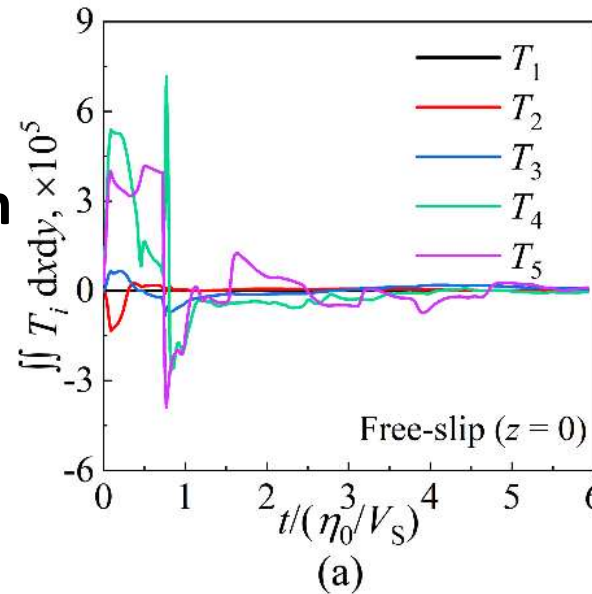
- Less vorticity near the wall
- Damping of local flame perturbation near the wall
- Non-uniform perturbation evolution





Vortex Dynamics

- T_1 Vortex stretching
- T_2 Viscous torque
- T_3 Viscous dissipation
- T_4 Dilation
- T_5 Baroclinic torque



$$\frac{\partial \omega_i}{\partial t} + u_k \frac{\partial \omega_i}{\partial x_k} = \underbrace{\omega_k \frac{\partial u_i}{\partial x_k}}_{T_1} - \underbrace{\varepsilon_{ijk} \frac{1}{\rho} \frac{\partial \rho}{\partial x_j} \left(\frac{\partial \tau_{kl}}{\partial x_l} \right)}_{T_2} + \underbrace{\frac{\varepsilon_{ijk}}{\rho} \left(\frac{\partial^2 \tau_{kl}}{\partial x_j \partial x_l} \right)}_{T_3} - \underbrace{\omega_i \frac{\partial u_k}{\partial x_k}}_{T_4} + \underbrace{\frac{1}{\rho^2} \varepsilon_{ijk} \frac{\partial \rho}{\partial x_j} \frac{\partial p}{\partial x_k}}_{T_5}$$



Conclusions

- Flame perturbation during shock-flame interaction can be divided into two stages, **shock compression** and **perturbation growth stages**.
- Wall friction has a significant influence on the shock-flame interaction and flame perturbation growth. Two effects of wall friction on flame-shock interaction:
 - flame stretching**
 - damping of local flame perturbation near the no-slip wall**
- **The wall friction promotes entire flame instability after SFI rather than weakening.**



Acknowledgment

**National Key Research and Development Program of China
(Grant No. 2021YFB4000902)**

**Fundamental Research Funds for the Central Universities
(Grant No. WK2320000055)**

DNL Cooperation Fund

China Scholarship Council



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Thank you for listening!

Q&A

